

Deep spectroscopy of 9C J1503+4528: a very young compact steep spectrum radio source at $z = 0.521$

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ABSTRACT

9C J1503+4528 is a very young compact steep spectrum radio galaxy, with an age of the order of 10^4 yr. This source is therefore an ideal laboratory for the study of the intrinsic host galaxy/intergalactic medium properties, interactions between the radio source and surrounding interstellar medium, links between star formation and active galactic nucleus (AGN) activity and the radio source triggering mechanism. Here we present the results of a spectroscopic analysis of this source, considering each of these aspects of radio source physics.

We find that shock ionization by the young radio source is important in the central regions of the galaxy on scales similar to that of the radio source itself, whilst evidence for an AGN ionization cone is observed at greater distances. Line and continuum features require the presence of a young stellar population (YSP), the best-fitting model for which implies an age of 5×10^6 yr, significantly older than the radio source.

Most interestingly, the relative sizes of radio source and extended emission-line region suggest that both AGN and radio source are triggered at approximately the same time.

If both the triggering of the radio source activity and the formation of the YSP had the same underlying cause, this source provides a sequence for the events surrounding the triggering process. We propose that the AGN activity in 9C J1503+4528 was caused by a relatively minor interaction, and that a supermassive black hole powering the radio jets must have been in place before the AGN was triggered.

Key words: galaxies: active – galaxies: evolution – galaxies: individual: 9C J1503+4528 – galaxies: ISM – galaxies: stellar content – radio continuum: galaxies.

1 INTRODUCTION

The triggering of powerful radio sources is still poorly understood. It is generally thought that some mechanism must channel significant quantities of gas into the central regions of the host to fuel the active galactic nucleus (AGN), wherein a rapidly spinning black hole may power the radio jets (e.g. Blandford & Znajek 1977).

Two scenarios have been proposed. The first is the infall of gas which has cooled in the halo of a galaxy, or a cooling flow in a cluster (Bremer, Fabian & Crawford 1997); this proposition has received much interest recently because it may have a thermostat

feedback effect to prevent runaway cooling in clusters (Brüggen & Kaiser 2002; Fabian et al. 2003). However, it is most likely that this mechanism in fact causes a renewal of activity in an existing radio source (Nipoti & Binney 2005), rather than being related to the original triggering.

The second is associated with the long-standing idea that a merger between two galaxies, at least one of which is gas rich, is a precursor to AGN activity (e.g. Heckman et al. 1986). Wilson & Colbert (1995) suggested that the coalescence of black holes following a major merger creates the necessary rapidly spinning black hole, which suggests links between radio sources and other merger systems such as ultraluminous infrared galaxies, and implies that evidence for young stellar populations (YSPs) and disturbed morphologies should be present at some level in observations of young radio sources.

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The merger scenario is supported by an increasing amount of anecdotal evidence. For instance, there are several cases known where the host galaxy of a powerful radio source appears to be undergoing a major merger (e.g. Johnston et al. 2005). Similarly, there is increasing evidence that powerful high-redshift radio galaxies lie in forming clusters (see e.g. Venemans et al. 2005); such environments would be rich in major mergers. This evidence in turn takes us back to the first scenario above; massive cluster galaxies may host a spinning black hole as a relic of their early history, which is generally quiescent but which becomes active when fuelled by cooling gas. It is clearly evident that radio sources play a significant part in hierarchical galaxy and structure formation (e.g. Croton et al. 2006), so the precise ingredients for radio source formation are a cause of renewed interest.

In any case, our understanding of this early phase of radio galaxy evolution is hampered by the difficulty of compiling uniform samples of radio sources in the first few millennia of their growth, and indeed the impossibility of pinpointing galaxies immediately prior to this stage. However, it is now widely accepted that compact steep spectrum (CSS) and some gigahertz-peaked spectrum (GPS) radio sources represent the earliest stages of radio source evolution, being the progenitors of the larger Fanaroff–Riley type II (FR II) source population (e.g. Fanti et al. 1995; Readhead et al. 1996; Snellen et al. 2000; Murgia 2003). The study of such sources in the stages immediately after triggering may be a valuable way of elucidating the events that have occurred beforehand.

In an effort to secure a uniform sample of young sources, we have turned to the 9C survey (Waldram et al. 2003). This was carried out at 15 GHz with the Ryle Telescope to identify foreground sources for the Very Small Array (VSA) microwave background experiment. Over the past several years, extensive follow-up of the radio properties of this population has been carried out (Bolton et al. 2003, 2004, 2006), both to understand the effects on microwave background experiments better and to investigate the sources themselves; the high survey frequency has proved a good way of selecting

significant numbers of CSS and GPS sources, as well as sources peaking in the ~ 10 GHz range. We are now undertaking optical followup of the CSS and GPS sources, and in this paper we present a detailed spectroscopic study of one of these compact sources, 9C J1503+4528, a classical double FR II radio galaxy at a redshift of $z \sim 0.5$ with an angular size of 0.5 arcsec.

Details of the observations are presented in Section 2, analysis of the emission-line strengths and ionization state in Section 3, gas kinematics in Section 4, emission morphology in Section 5 and a discussion of the implications of our results in Section 6. We give a summary of our conclusions in Section 7.

Values for the cosmological parameters of $\Omega_0 = 0.27$, $\Omega_\Lambda = 0.73$ and $H_0 = 65 \text{ km s}^{-1} \text{ Mpc}^{-1}$ are assumed throughout.

2 THE OBSERVATIONS

The observed source, 9C1503+4528 (see Fig. 1), was selected as our first target for detailed study for several reasons.

First, the source has a rest-frame 1.4-GHz luminosity of $1.33 \times 10^{25} \text{ W Hz sr}^{-1}$, well above the FR I/FR II break. Secondly, the source is resolved, and clearly shows a classical double structure in the radio map presented in Bolton et al. (2004). The angular size of the source is 0.5 arcsec, corresponding to a projected linear size of $\approx 3 \text{ kpc}$. It is currently accepted that GPS and CSS sources are at a young stage in the evolution of the radio source (see e.g. Snellen et al. 2000). The picture in which sources are young is strongly supported by observations of the hotspot advance speed in many of the most compact sources. These observations find hotspot advance speeds in the range $0.1\text{--}0.4c$ (e.g. Owsianik, Conway & Polatidis 1998; Taylor et al. 2000; Polatidis & Conway 2003), which in the case of 9C J1503+4528 implies that the source has an age of the order of $\sim 10^4 \text{ yr}$. This places the source very close to its time of triggering, since the typical lifetime of an FR II source is $\sim 10^7$ to 10^8 yr (e.g. Kaiser, Dennett-Thorpe & Alexander 1997). Finally, the source redshift ($z \sim 0.5$) makes many useful diagnostic lines and the

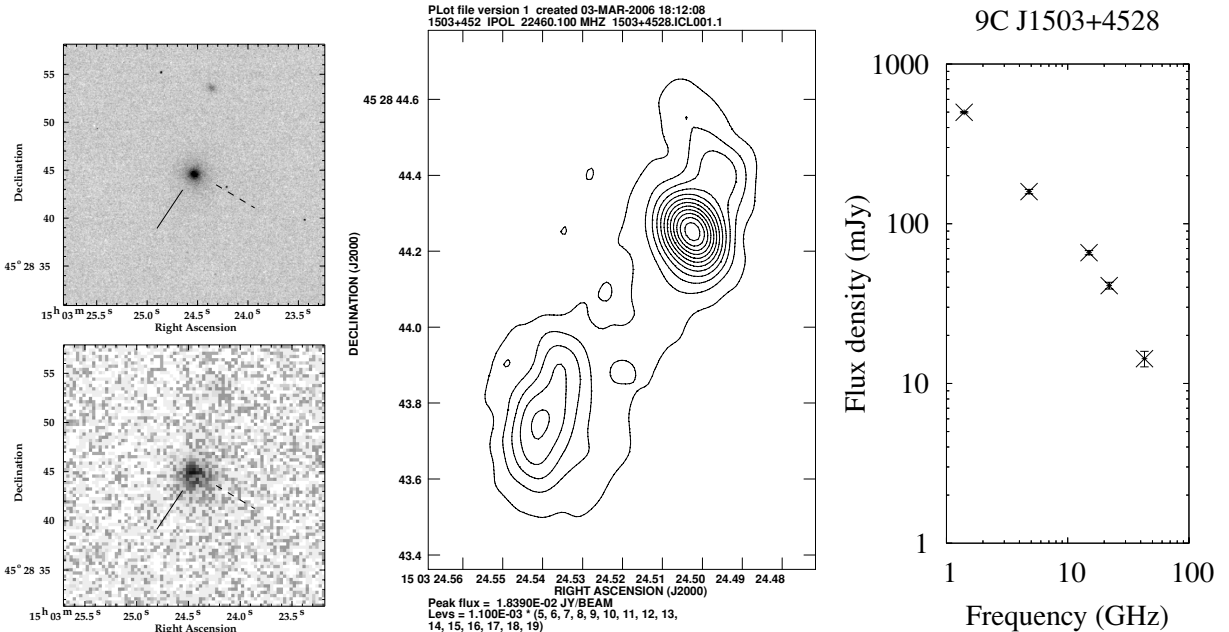


Figure 1. The compact radio source 9C J1503+4528. Left-hand panel: UFTI *K*-band (top panel) and WFC *r*-band (bottom panel) images of the host galaxy, an $R \approx 19$ elliptical. The two spectroscopic slit PAs, parallel and perpendicular to the radio source axis are indicated on each diagram by the solid and dashed lines, respectively. Middle panel: VLA 22-GHz radio contours. Right-hand panel: VLA and OVRO radio spectrum for this source, from Bolton et al. (2004).

4000-Å break accessible to moderate-resolution spectroscopy in a wavelength range region relatively clear of night-sky emission. With these facts in mind, the young CSS radio source 9C J1503+4528 is therefore a particularly appropriate object for the study of the intrinsic host galaxy/intergalactic medium (IGM) properties, interactions between the radio source and surrounding interstellar medium (ISM), links between star formation and AGN activity and the underlying radio source triggering mechanism.

2.1 Very Large Array imaging

The 22-GHz data obtained by Bolton et al. (2004) were re-analysed to obtain the map shown in Fig. 1. Both the central component and the SE compact component are unresolved. The most obvious interpretation of the radio morphology is that the bright central component is the core, with the SE compact component being a hotspot, and the SE and NW extensions being lobes. Given the overall steep spectrum, it is implausible that the source is being viewed at an angle close to the line of sight; in such a case, a strongly boosted flat-spectrum jet component would be expected to dominate in this high-frequency map. We thus conclude that the source has a double-lobed structure with an overall extent of ~ 3 kpc.

2.2 Optical and infrared imaging

An image using the Isaac Newton Telescope Wide Field Camera *r*-band filter was taken on the night of 2005 April 11. This filter has the advantage of not including the strong [O II] and [O III] emission lines at the source redshift. Five 600-s images were taken, offset at the centre and corners of a 20-arcsec pattern, and stacked. The resulting image is shown in Fig. 1.

A *K*-band image was also obtained, on the night of 2006 March 5 using the UFTI instrument at UKIRT. Nine 60-s jittered observations were taken and reduced using the standard ORAC-DR pipeline; the resulting image is also displayed in Fig. 1. The seeing, measured from an observation of UKIRT standard star P177-D taken immediately after the target observation, was 0.6 arcsec.

2.3 Optical spectroscopy

The spectroscopic observations were carried out on 2004 July 18, using the Keck LRIS spectrograph (Oke et al. 1995). Seeing was ~ 0.8 arcsec, and a 1-arcsec slit was used. The 600/7500 grating was used, providing a wavelength range of ~ 5200 – 7800 Å, a spatial scale of 0.215 arcsec pixel $^{-1}$, a spectral scale of 1.26 Å pixel $^{-1}$ and a spectral resolution of ~ 4.4 Å. For an object lying at a redshift of $z \sim 0.5$, the corresponding rest-frame spectral range is approximately 3500 – 5200 Å, thereby including emission lines from [O II] 3727 Å up to [O III] 5007 Å. Slit position angles (PAs) of $+45^\circ$ and -45° were chosen, aligning the slit either parallel or perpendicular to the radio source axis. Long-slit spectra were obtained with total integration times of 1800s at each slit PA. The integrations were split into 900-s exposures to aid the removal of cosmic rays.

Standard packages within the NOAOIRAF reduction software were used to reduce the raw data. Corrections were made for over-scan bias subtraction, and the data were then flat-fielded using internal lamps. The flat-fields were taken immediately after each observation of the target, to minimize the effects of fringe drift in the CCD. The two-dimensional frames were wavelength calibrated and corrected for distortions in the spatial direction using wavelength-calibration arc exposures. Taking care not to subtract any of the extended line emission, the sky background was removed. Observations of the smooth-spectrum standard star BD +28 4211 were

used to provide accurate flux calibration of the final spectra and to remove the telluric *A*- and *B*-band absorption features. The central regions of the parallel and perpendicular spectra effectively probe the same region of the host galaxy. These two spectra were combined in order to produce a higher signal-to-noise ratio (S/N) spectrum, for use in analysing the emission from the central regions of the radio galaxy.

The resulting parallel and perpendicular two-dimensional spectra are displayed in Fig. 2. In the parallel spectrum, the presence of narrow emission, extended over a distance of ~ 5 arcsec (approximately 10 times larger than the radio source itself) can clearly be seen, in addition to a broader component lying close to the continuum centroid. The [O II] emission is particularly strong compared to [O III], and the continuum emission short of the 4000-Å break is also substantially more luminous than would be expected for an old stellar population alone. The perpendicular spectrum displays similar properties, with the exception that extended narrow emission is not observed. A one-dimensional spectrum extracted from the central 2 arcsec of the combined perpendicular and parallel spectra is displayed in Fig. 3.

Determining the precise systemic redshift of 9C J1503+4528 was complicated by the inherent structure in the emission lines: the redshift implied by the peak of the line emission on the continuum centroid differs substantially from that of the narrow extended component. However, the various absorption features in the spectra (Ca H & K, high-order Balmer lines) all suggest a redshift identical to that determined from the extended narrow component, giving a value of $z = 0.521$.

3 EMISSION-LINE FLUXES AND THE IONIZATION STATE OF THE GAS

One-dimensional spectra were extracted from the parallel, perpendicular and combined two-dimensional frames. The extractions were carried out over 2 arcsec (~ 12 kpc) for the parallel, perpendicular and combined frames, and additionally over 5 arcsec (~ 31 kpc) for the parallel and perpendicular frames. The lines were then fitted by Gaussians; the resulting fluxes and equivalent widths are presented in Appendix A.

Spectra were also extracted in 0.5-arcsec steps across the slit, in order to investigate the changing ionization state of the emission-line gas with position via ionization state diagnostic diagrams (Baldwin, Phillips & Terlevich 1981). For an effective analysis, model predictions should be well separated and results should be consistent over several different diagrams/line combinations. Line pairs should ideally be close in wavelength and of the same species, to limit the effects of reddening and variations in gas composition. The most appropriate lines available from our optical spectra are [O II]3727 Å, [O III]4363 Å, H β and [O III]5007 Å.

Fig. 4 displays the diagnostic diagram for the line pairs [O II]3727 Å/H β and [O III]5007 Å/H β , extracted from the parallel spectra in 0.5-arcsec steps, together with various theoretical model tracks. These include: (i) the shock models of Dopita & Sutherland (1996), which also incorporate the effects of a precursor ionization field, (ii) the mixed medium (matter bounded versus ionization bounded) photoionization tracks of Binette, Wilson & Storchi-Bergmann (1996), plus (iii) simple AGN photoionization tracks (Groves, Dopita & Sutherland 2004a,b), also incorporating the effects of dust. The photoionization models displayed are for solar metallicity and a density of $n = 1000$ cm $^{-3}$. For reference, the photoionization tracks display little variation for lower values of n , while the $n = 10000$ cm $^{-3}$

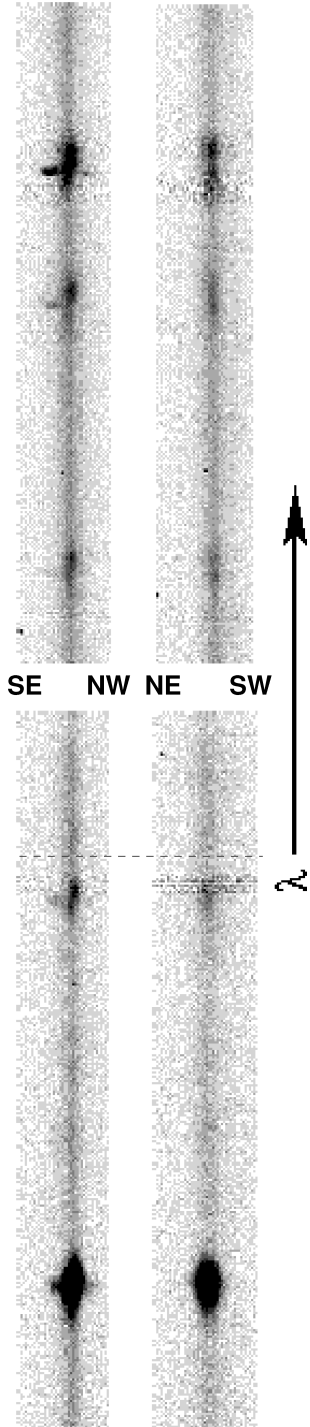


Figure 2. Sections of particular interest (rest frame 3730–4070 and 5190–5490 Å) from the two-dimensional optical spectra of 9C1503+4528, with the slit aligned parallel (left-hand side) and perpendicular (right-hand side) to the radio source axis. Wavelength increases from the bottom to the top of the diagram, as indicated by the direction of the arrow, and the position of the 4000-Å break is denoted by the dashed line. The extracted regions are approximately 10-arcsec wide. The five prominent emission lines are (bottom to top): [O II]3727, [Ne III]3869, H β , [O III]4959 and [O III]5007 Å.

track is shifted to smaller values of $\log([\text{O II}]/\text{H}\beta)$ by ~ 0.2 dex. The 2- Z_{\odot} metallicity tracks are generally similar to those of solar metallicity, whereas lower Z tracks are shifted to smaller values of both $\log([\text{O II}]/\text{H}\beta)$ and $\log([\text{O III}]/\text{H}\beta)$.

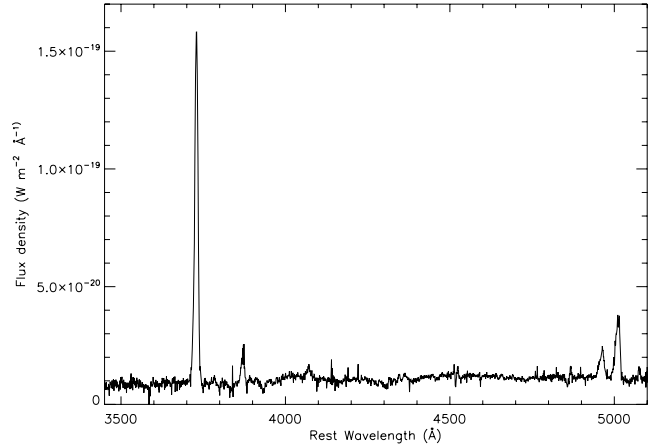


Figure 3. Spectrum of 9C J1503+4528 obtained from the central 2 arcsec of the combined parallel and perpendicular data. This spectrum provides the highest S/N for our data, and provides a good illustration of the relative emission line strengths and the fairly blue colour of the continuum.

This diagram clearly shows the changing ionization state relative to distance from the nucleus. The spectra extracted from the central regions of the host galaxy lie close to the predictions of the shock ionization models (and also the far end of the mixed medium ionization track). At intermediate distances (~ 1 arcsec, i.e. sampling emission from just outside the radio source), the shock plus precursor photoionization models are most suited. Beyond 1–1.5 arcsec (i.e. the extended narrow emission region), the data are well described by simple AGN photoionization. High values of the ionization parameter, U , are the most plausible, and a dusty medium is favoured. For the photoionized gas, the ionization state at larger radii suggests either a lower spectral index ($\alpha \sim -2.0$ rather than $\alpha \sim -1.2$), higher gas density or lower metallicity. The mixed medium models of Binette et al. only provide a good explanation for a few of the data points, and it seems unlikely that this system can be well described using that particular model. The observed line ratios are not likely to be affected by internal reddening, as the Balmer line ratios in our spectra are in good agreement with the predictions of case B recombination. Balmer line absorption could potentially have some influence on the exact positioning of the data points in this figure; while we believe we have correctly accounted for the errors on our measurement of H- β , additional shifts of up to 0.2 dex towards smaller line ratios cannot be completely ruled out. However, this would not have any effect on our overall interpretation of the data.

The picture provided by these results is as follows. The ultraviolet (UV) continuum emission from the AGN ionizes the gas at all locations within the ionization cone, decreasing in intensity at larger distances. The innermost regions of the galaxy, of the order of the same size as the fledgling radio source, have an additional ionization contribution from shocks. The UV radiation field produced by these shocks adds another ionization component to the emission from gas at intermediate distances.

The relative ionization states of the parallel and perpendicular spectra can be compared in Fig. 5. This figure displays the same tracks as in Fig. 4, for spectra extracted from various regions of our two-dimensional spectra. These include the central 2-arcsec regions of the combined, parallel and perpendicular spectra. The results for the central 1–2 arcsec should not be greatly different, as both spectra sample essentially the same region. However, the impact of AGN

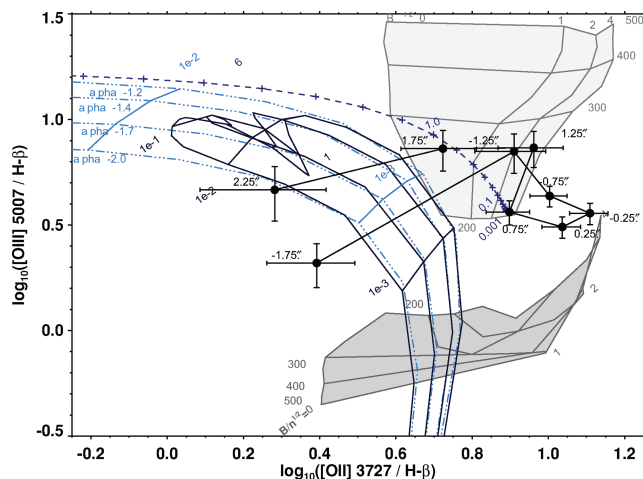


Figure 4. Ionization mechanism diagnostic plot for 9C1503+4528, using the line ratios $[\text{O II}]\lambda 3727/\text{H}\beta$ and $[\text{O III}]\lambda 5007/\text{H}\beta$. Data points represent spectra extracted in 0.5-arcsec steps, and are labelled in terms of their distance (in arcsec) from the continuum centroid. The model tracks include: (1) shock ionization (Dopita & Sutherland 1996), with and without a precursor photoionization region (light and dark shading, respectively), (2) the matter-bounded photoionization track (dashed dark blue track, Binette et al. 1996) and (3) simple AGN photoionization tracks (light blue tracks), including the effects of dust (solid black tracks; Groves et al. 2004a,b). The photoionization models are for $n = 1000 \text{ cm}^{-3}$, solar metallicity and various values of the ionization index α (from $1.2 < \alpha < 2.0$). A brief description of the effects of changes in metallicity and/or gas density is given in the main text; we refer the reader to Groves et al. (2004a,b) for full details.

and shock ionization will be greatly lessened at larger distances, for material which lies off the radio source axis and well outside the ionization cone of the AGN. Also displayed are the observed emission line ratios for spectra extracted at distances from the continuum centroid of $+1.5$ to $+2.5$ arcsec and -1.5 to -2.5 arcsec, for both the parallel and perpendicular spectra. The large errors on the data points for the extended regions of the perpendicular spectra reflect

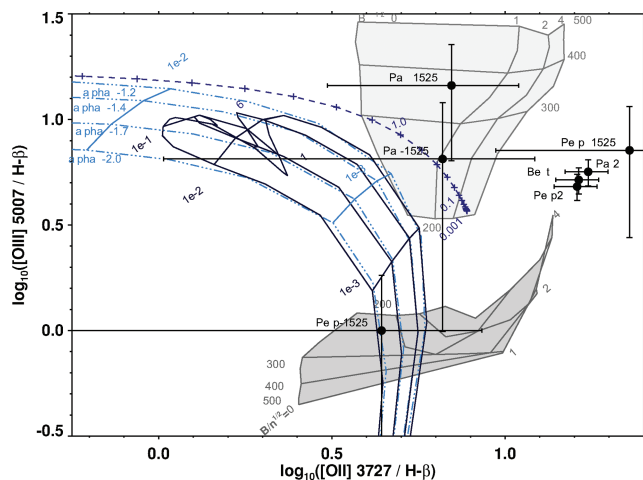


Figure 5. Ionization mechanism diagnostic plots for 9C1503+4528, using the line ratios $[\text{O II}]3727/\text{H}\beta$ and $[\text{O III}]5007/\text{H}\beta$. Data points represent spectra extracted from various regions of the spectra taken perpendicular and parallel to the radio axis, and include: (i) the central 2 arcsec of the parallel, perpendicular and combined (best S/N) spectra, (ii) 1-arcsec wide sections of the parallel and perpendicular spectra lying ± 2 arcsec from the continuum centroid. The model tracks are the same as those in Fig. 4(a).

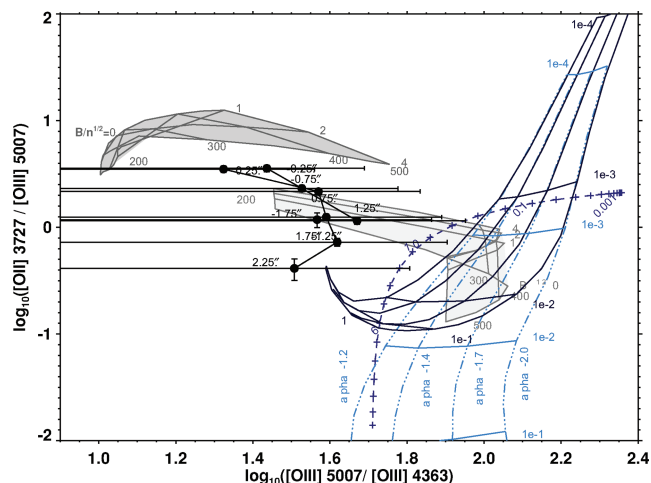


Figure 6. Ionization mechanism diagnostic plots for 9C1503+4528, using the line ratios $[\text{O II}]\lambda 3727/[\text{O III}]\lambda 5007$ and $[\text{O III}]\lambda 5007/[\text{O III}]\lambda 4363$. Data points represent spectra extracted in 0.5-arcsec steps, and are labelled in terms of their distance from the continuum centroid. The model tracks are the same as those in Fig. 4(a), and include: (1) shock ionization (Dopita & Sutherland 1996), with and without a precursor photoionization region (light and dark shading, respectively), (2) the matter-bounded photoionization track (dashed dark blue track, Binette et al. 1996) and (3) simple AGN photoionization tracks (light blue, various styles), including the effects of dust (solid black tracks, Groves et al. 2004a,b).

the very weak line emission present at these distances. In the case of the -1.5 to -2.5 arcsec data point, shocks are not a plausible option, and the data should lie outside the ionization cone of the AGN. Instead, the position of this data point suggests that ionization by young, hot stars may provide a better explanation (see e.g. Kewley, Geller & Jansen 2004 and Rola, Terlevich & Terlevich 1997 for appropriate locations on the $[\text{O II}]/\text{H}\beta$ versus $[\text{O III}]/\text{H}\beta$ diagnostic plot for H II regions; such emission line ratios would generally be expected to lie towards the lower left-hand quadrant of Fig. 5). A combination of ionization by young stars and the AGN could also be relevant for the -1.75 and $+2.25$ arcsec data points on Fig. 4.

Fig. 6 shows the diagnostic diagram for the $[\text{O III}]5007\text{ \AA}/[\text{O III}]4363\text{ \AA}$ and $[\text{O II}]3727\text{ \AA}/[\text{O III}]5007\text{ \AA}$ line ratios, together with the observed data from our parallel spectra in spatial steps of 0.5 arcsec . Note that the mixed medium model track, which bent towards the data points at low values of $A_{M//1}$ in Fig. 4, tends towards higher values of $[\text{O III}]5007\text{ \AA}/[\text{O III}]4363\text{ \AA}$ on Fig. 6 as opposed to the lower values of this line ratio observed in the data. Although the error bars are large for the $[\text{O III}]5007\text{ \AA}/[\text{O III}]4363\text{ \AA}$ line ratio, we nevertheless have a certain amount of consistency with the previous diagnostic diagrams: once again the central regions of the galaxy favour the predictions of shocks, while more distant locations lie closer to the dusty AGN photoionization tracks. The relatively low $[\text{O III}]5007\text{ \AA}/[\text{O III}]4363\text{ \AA}$ ratio is indicative of high temperatures and/or densities in the emitting gas, of the order of several 10^4 K in the low-density limit (e.g. Osterbrock 1989), fully consistent with the presence of shocks in these regions.

¹The solid angle ratio of matter- to ionization-bounded material in the mixed medium model of Binette et al. (1996).

4 GAS KINEMATICS

We have also investigated the changing gas kinematics, via an analysis of the profiles of the [O II] 3727 Å and [O III] (4959+5007 Å) lines. Two-dimensional regions around each of these emission lines were extracted, and from these a sequence of one-dimensional spectra were extracted along the slit direction, stepped every 0.5 arcsec.

After allowing for continuum subtraction, the data were fitted by up to three successive Gaussian components, which were only accepted if their full width at half-maximum (FWHM) was larger than the instrumental resolution and the S/N was greater than approximately five (excepting the outermost spectra, for which a fainter, single Gaussian component could be detected at a lower S/N). The combination of Gaussians with the lowest reduced χ^2 were then taken to be the maximum number of velocity components which could be fitted to the data (we also fixed the relative positions and strengths of the 4959- and 5007-Å lines). Example fits are displayed in Fig. 7.

The integrated line flux, the velocity relative to that at the centre of the galaxy, and the emission-line FWHM were determined for each Gaussian fit. We calculated the FWHM by subtracting in quadrature the instrumental FWHM, as determined from unblended skylines. We then calculated errors for these three parameters, allowing for the fact that ranges of possible fits are equally acceptable at low S/N. This approach allows us to search for high-velocity components in the emission-line gas, or other structures incompatible with a fit to a single-velocity component. Given the observed emission line profile of this source, such a step was a necessity, and broader emission-line

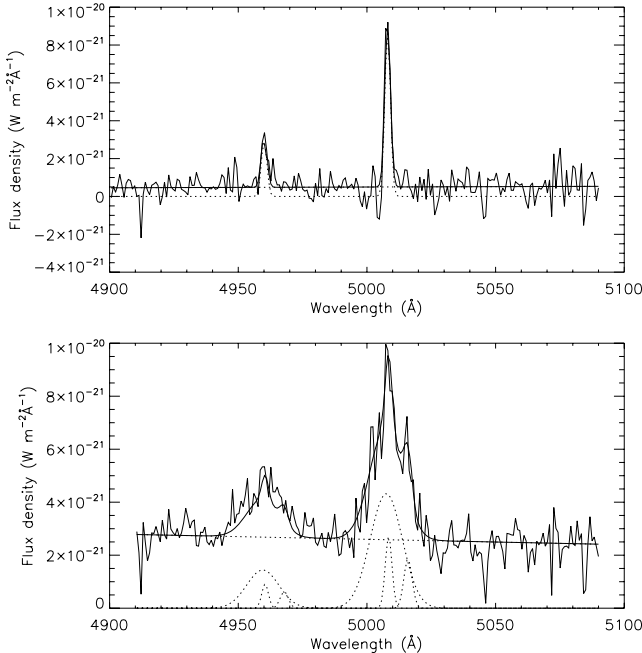


Figure 7. Example of line profile fits for the [O III](4959+5007) Å emission lines. The top figure is for the spectrum extracted at a distance of 2.25 arcsec to the SE of the continuum centroid, and the bottom figure for a region lying at a distance of 0.75 arcsec in the same direction. The data were modelled with several components (dotted lines): a continuum level, and one to three separate Gaussian components for each line. For [O III], the relative strengths and positions of the 4959- and 5007-Å lines were fixed. The combination of Gaussians with the lowest reduced χ^2 determined the maximum number of velocity components which could be fitted to the data.

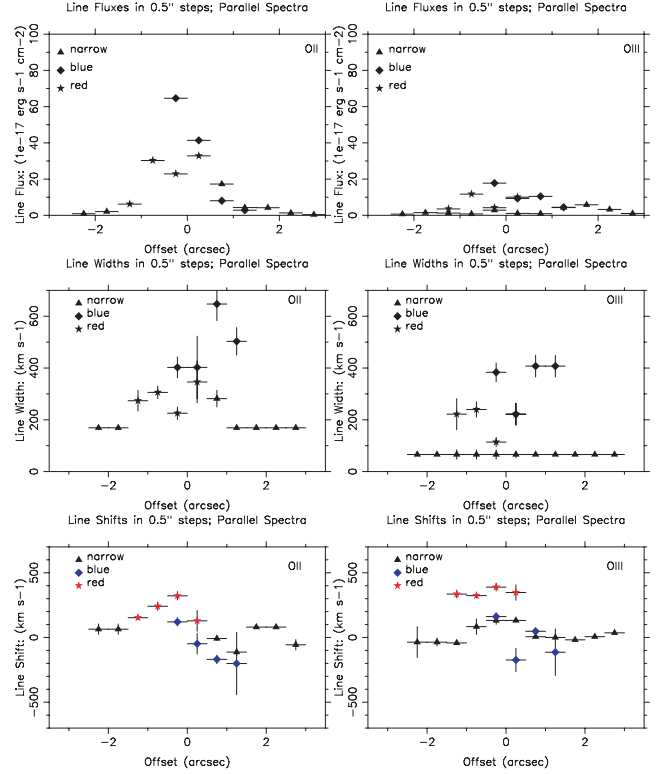


Figure 8. Gas kinematics for 9C J1503+4528. The changing line fluxes (top panel), widths (middle panel) and offsets (bottom panel) with position are displayed for the [O II] (left-hand panel) and [O III] (right-hand panel) emission lines, extracted and fitted in 0.5 steps from the parallel spectrum. Up to three emission-line components are modelled at each position. Negative offsets are towards the NW, and positive offsets towards the SE.

components are clearly present in the central parts of the extended emission-line region (EELR).

The results of our analysis of the kinematics of the parallel spectrum are presented in Fig. 8.

An unresolved emission component is present in the outer regions of the EELR, in both [O II] and [O III], at wavelengths consistent with the emitting material lying at rest relative to the host galaxy. This narrow component can be traced across the galaxy in [O III]. The [O II] emission line is of substantially higher luminosity, particularly in the central regions of the galaxy. Because of this, the narrow component of [O II] can only be observed in the outer regions, and is lost in the centre in favour of one or more broader components. The relative strengths of the narrow components of [O II] and [O III] are consistent with simple AGN photoionization, as demonstrated by our diagnostic diagrams in the previous section.

The broader emission-line components extend to approximately 1–1.5 arcsec from the continuum centroid. Note that the [O II] emission is substantially stronger than the [O III] emission; such a low-ionization state in the emitting material is suggestive of shocks associated with the radio source, in which case the emitting material would display larger linewidths than more settled material. Usually a single broad component is sufficient, except in the central regions where two broad components are preferred. The broad components on the SE side of the galaxy (see the two-dimensional spectrum presented in Fig. 2) usually lie at shorter wavelengths, with a velocity shift of a few hundred km s^{−1}, whilst the emission from the opposite side of the galaxy is shifted to longer wavelengths by a similar

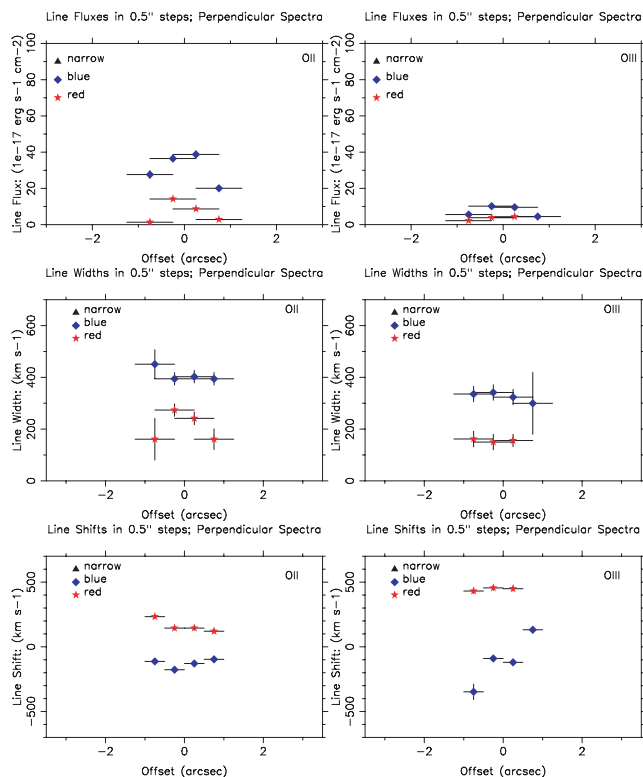


Figure 9. Gas kinematics for 9C J1503+4528. The changing line fluxes (top panel), widths (middle panel) and offsets (bottom panel) with position are displayed for the [O II] (left-hand panel) and [O III] (right-hand panel) emission lines, extracted and fitted in 0.5-arcsec steps from the perpendicular spectrum. Up to three emission-line components are modelled at each position. Negative offsets are towards the SW, and positive offsets towards the NE.

amount. Of these two distinct components, the blueshifted emission-line component is generally more luminous, and has a larger linewidth (400–600 km s^{−1}, cf. 200–400 km s^{−1} for the redshifted emission).

We also present the results of a similar analysis of the spectrum aligned perpendicular to the radio source axis (Fig. 9). As the line emission in the perpendicular spectrum is much less extensive, larger extracted spectra of 1 arcsec in width have been modelled to boost the S/N. The first main feature of note is the lack of an extended narrow component. This is strong evidence that the source of the ionizing photons for the extended narrow component is the ionization cone of the AGN, which does not intersect with the perpendicular spectrum except in its centre.

In the central regions of the perpendicular spectrum, broad-emission components are observed. These are very similar to those in the parallel spectrum, with the exception that there is no clear tendency for the blue or redshifted emission to preferentially lie on one side of the galaxy or the other. This suggests that outflows along the radio axis are a plausible option for explaining the gas kinematics, as an alternative to simple rotation. The greater luminosity and extent of the material to the SE side of the galaxy, together with the blueshifted broader component, suggests that the SE radio lobe would be orientated towards the observer; given the observed radio asymmetries (Section 2.1) and quiescent extended emission-line component at larger radii the outflow scenario is indeed highly plausible.

5 CONTINUUM EMISSION

The continuum emission of 9C1503+4528 is remarkably blue. Together with the clear presence of Balmer absorption lines and the young age of the radio source this suggests a significant YSP. If the triggering of radio source activity is connected with the star formation, the properties of any YSP formed as part of that process can shed light on the exact mechanisms and time-scales involved. We have therefore carried out continuum modelling of our spectra, with the aim of placing limits on both the mass and age of any YSP present in this galaxy.

For the continuum modelling, we combined both the parallel and perpendicular two-dimensional frames and extracted a one-dimensional spectrum from the central 2 arcsec (Fig. 3).

The first stage was to remove the contribution of nebular continuum emission (Dickson et al. 1995). Although subject to stellar absorption features, the Balmer line strengths are in general consistent with the predictions of case B recombination, and internal reddening is unlikely to be a major issue. The H β flux can be well modelled, and was used to determine the contribution from the various nebular continuum emission processes (free–free emission, bound–free emission, two-photon emission and the Balmer forest pseudo-continuum), assuming an electron temperature of 15 000 K. (Checks with alternative temperatures, allowing for the error on our evaluation of the H β flux, do not produce significantly different results.) The corrected spectrum was binned in blocks of ~ 30 Å, excluding any regions where emission lines were present.

To fit the remaining stellar continuum, we used a combination of young and old stellar population spectral energy distributions (SEDs) from the models of Bruzual & Charlot (2003). A Chabrier initial mass function and solar metallicity were selected, and the stellar populations were assumed to form in an instantaneous burst. We consider stellar population ages ranging from $<10^5$ up to 9×10^9 yr (the expected age of the old stellar population in a $z \sim 0.5$ galaxy, assuming a formation redshift of $z \sim 10$ –20). Reduced- χ^2 fitting was used to determine the relative proportions of old/young stars for all combinations of SED models.

Fig. 10 displays the results of our best-fitting model. This consists of an old (6 Gyr) stellar population, plus a 5×10^6 yr YSP accounting for 0.065 per cent of the total stellar mass, and roughly 70 per cent of the total continuum flux at a wavelength of ~ 3600 Å. Our modelling implies that the total stellar mass in the central region of the galaxy probed by our spectra is approximately $10^{10} M_\odot$, whereas the overall stellar mass of the galaxy (based on our K -band modelling in Section 6) is closer to $10^{11} M_\odot$. The largest residuals lie near the Ca absorption feature at ~ 3970 Å, which is subject to infilling by line emission at 3967 Å due to H ϵ + [Ne III].

A plot of best-fitting reduced χ^2 as a function of YSP age is displayed in Fig. 11. A contribution from a YSP is very much favoured over an old stellar population alone. Fits obtained using different ages for the majority old stellar population (ranging from 5 to 9 Gyr in age) do not lead to any significant variation in the resulting best-fitting YSP.

We must also consider the potential presence of a power-law ionizing continuum associated with the AGN. If the presence of a power-law component is simply added to our modelling, allowing its flux and shape to vary freely, the resulting best-fitting model generally prefers an unrealistically strong power-law flux, and it is therefore necessary to place constraints on the power-law contribution in order to make a meaningful assessment of its effects. To determine the balance between YSP age/mass and power-law contribution, we repeated our fitting procedure using fixed power-law

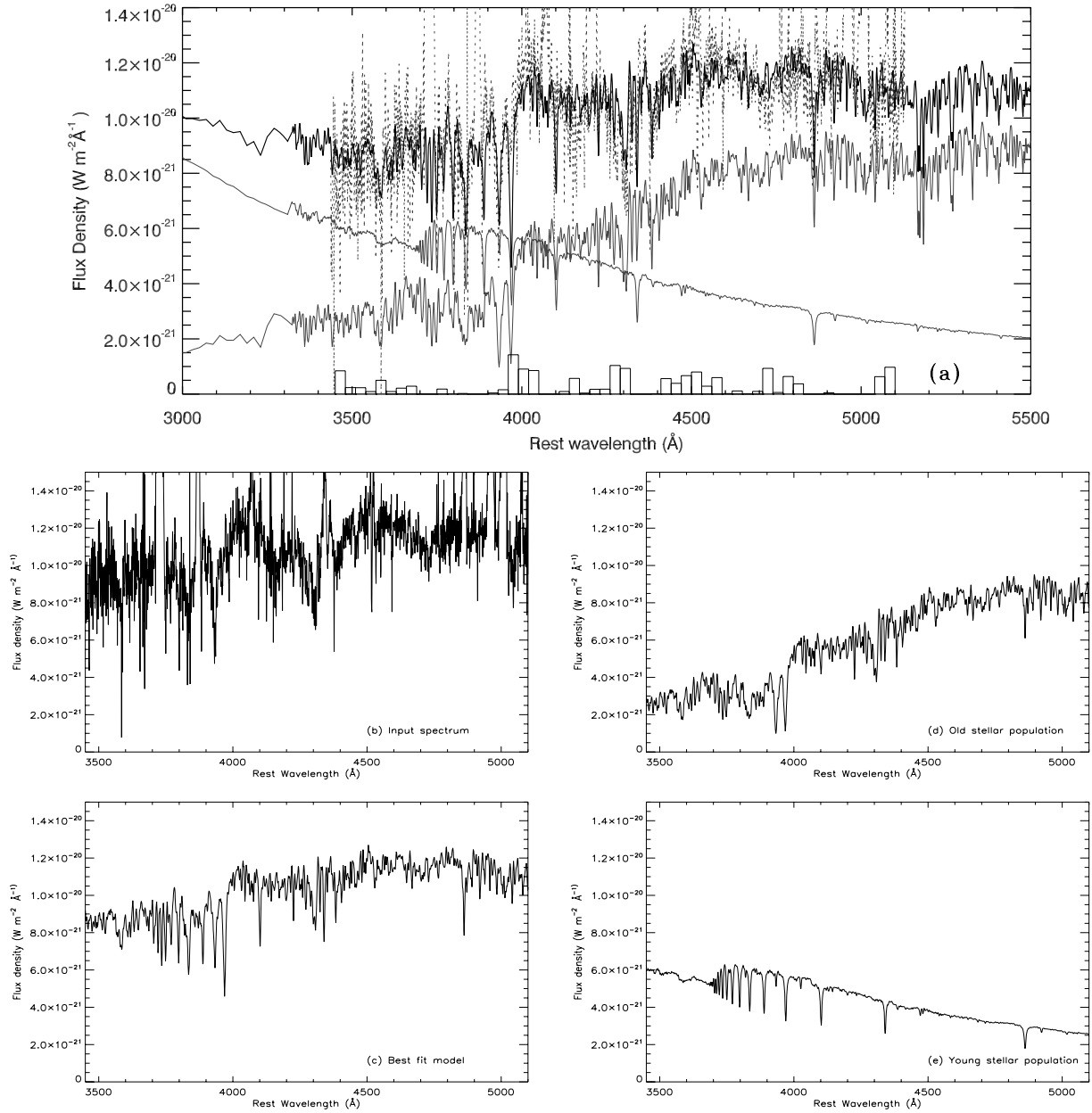


Figure 10. Continuum fits to the central 2 arcsec of the combined spectrum (Fig. 3). The top frame (a) displays the observed spectrum (dotted line), the best-fitting model (dark solid) and the two-component (old and young) SEDs which make up the best-fitting model. Also displayed in frame (a) are the model residuals for each of the bins used in the fitting procedure. For clarity, we also display the observed data [frame (b), centre left-hand panel], the best-fitting model [frame (c), lower left-hand panel], the old stellar population component [a 6-Gyr old stellar population, frame (d), centre right-hand panel] and the YSP component [a 5-Myr old stellar population, frame (e), lower right-hand panel] as separate spectra.

contributions ranging from 0 to 50 per cent of the total flux at a wavelength of $\lambda \sim 3600$ Å. The spectral index of the power law (defined as $F_\lambda \propto \lambda^\alpha$) is constrained to be one of three values, $\alpha = -1, -2$ or -3 . The fitting procedure was then repeated for all combinations of stellar populations and power-law parameters. The resulting best-fitting YSP parameters (age, mass and reduced χ^2) for each of the different power-law contributions considered are displayed in Appendix B.

With the exception of the fits with the very highest power-law contribution (> 25 per cent), the best-fitting YSP mass and age remain relatively unchanged. The similar strengths of the narrow line-emission components on either side of the galaxy centroid in the

parallel spectra (suggesting that this source is observed with the ionization cone aligned in a direction nearly perpendicular to the line of sight) also imply that a two-component stellar population provides a more than adequate explanation of the observed continuum emission, and that no power law is required.

6 IMPLICATIONS FOR SOURCE TRIGGERING

6.1 Morphological properties of the host galaxy

Given the spectroscopic evidence for a YSP, we examined the galaxy morphology for evidence for a merger or other interaction which

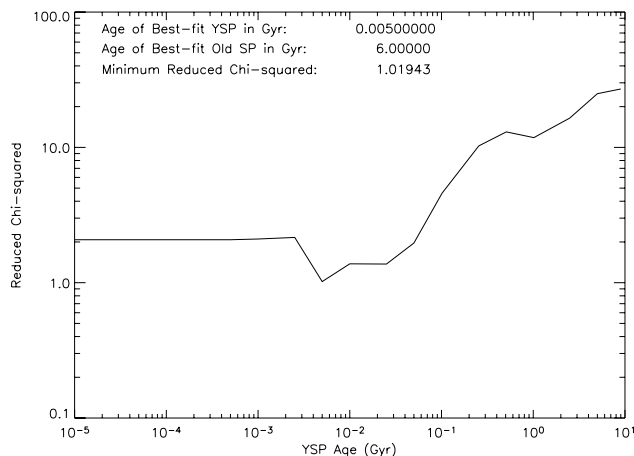


Figure 11. Reduced χ^2 of best-fitting model versus YSP component age for the continuum fitting of 9C J1503+4528. The best-fitting model includes a YSP component with an age of 5×10^6 yr, accounting for 0.065 per cent of the total stellar mass of the host galaxy, and roughly 70 per cent of the observed continuum flux at a wavelength of ~ 3600 Å. The age of the majority old stellar population in the best-fitting model is approximately 6 Gyr.

may have caused the starburst. Our calculated 0.06 per cent 5-Myr YSP would be expected to contribute only 0.02 mag to the *K*-band flux, and so the *K*-band image will be dominated by the old stellar populations of the host galaxy and any merging companion or tidal tails.

We fitted the galaxy morphology using the methods detailed in Inskip et al. (2005), which can be briefly summarized as follows. The point spread function for the UFTI field is determined from unsaturated stars within the mosaic field; we extract their two-dimensional profile, normalize to unit flux and take the average profile. Sérsic profile (Sérsic 1968) galaxy models were then convolved with the point spread function and fitted to the surface profile of the galaxy, using available least-squares minimization IDL routines.² Free parameters are the galaxy flux, centroid, effective radius, Sérsic index and fractional nuclear point-source contribution. Galaxy ellipticity was also allowed to vary, but in the case of this object a circular model with $\epsilon = 0$ was preferred.

From the *K*-band data, our resulting best-fitting model, residuals and one-dimensional radial profile are displayed in Fig. 12. We find that the host galaxy appears to have a simple de Vaucouleurs elliptical morphology, with a preferred Sérsic index of 4.0 ± 0.2 , and a best-fitting ($\chi^2 = 1.03$) effective radius of $r_{\text{eff}} = 1.70 \pm 0.15$ arcsec (equivalent to 11.5 ± 1.0 kpc at $z = 0.521$). We find no evidence for any nuclear point-source contribution at *K* band (preferred point-source percentage is < 1 per cent).

The *r*-band data are too noisy to make an accurate fit to the data; the effective radius is constrained only to lie between 1.3 ± 0.5 and 1.6 ± 0.5 arcsec. This is, however, consistent with the *K*-band fit. Most importantly, we find no significant residual emission which would indicate the presence of a close companion galaxy or tidal tails.

Any galaxy interaction/merger activity must therefore either have been very minor in nature, or have occurred sufficiently long ago that the old stellar population of the host galaxy has returned to

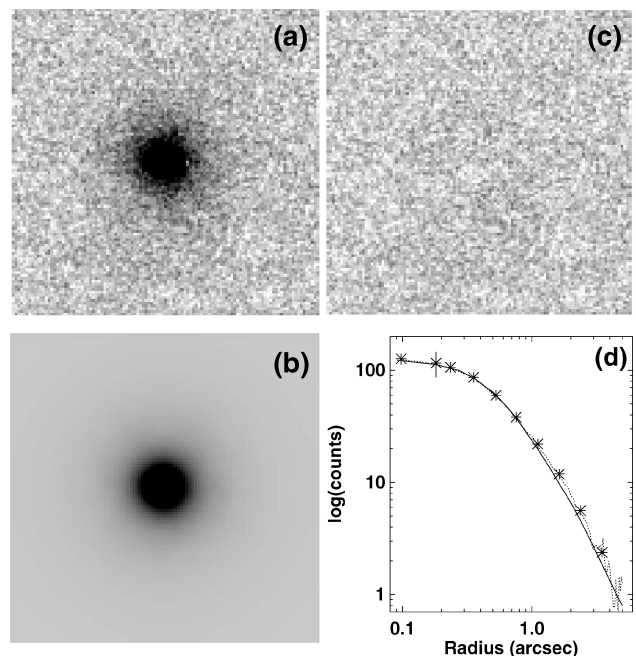


Figure 12. Host galaxy morphology fits for 9C J1503+4528. The *K*-band galaxy image is presented in frame (a), with the model galaxy given in frame (b) and the residuals after subtracting the best-fitting model from the data displayed in frame (c). Frame (d) illustrates a one-dimensional cut through the data (dotted line and binned points) with the best-fitting model (solid line) overlaid.

an undisturbed morphology. In the latter case, the interaction must have taken place several dynamical time-scales before the present observations. We estimate the dynamical time-scale T_{dyn} as r_{eff}/σ , where σ is the velocity dispersion of the galaxy, which in turn we estimate by taking the FWHM of the narrow component of the emission lines (~ 200 km s $^{-1}$). This gives $T_{\text{dyn}} \sim 60$ Myr, so any previous interaction must have occurred at a time very much earlier than the formation of the YSP. We thus conclude that the YSP and present AGN activity were caused by, at most, a minor interaction.

6.2 Evidence from the extended emission-line region

As with many other radio galaxies the properties of the EELR of 9C J1503+4528 vary substantially with spatial position. The changes in ionization state confirm our expectations that shocks are important in the vicinity of the young radio source, and that AGN photoionization dominates elsewhere. The kinematic properties of the emission-line gas confirm this picture, with the broader emission lines highlighting the disturbed kinematics of the central regions, in contrast to the kinematically quiescent material at larger radii.

The projected physical size of the EELR is ≈ 34 kpc. Emission-line regions of this scale are by no means unusual, and are often substantially larger in the case of FRII radio sources (for example, ~ 50 – 100 kpc for extended 3C and 6C sources, Best, Röttgering & Longair 2000; Inskip et al. 2002). However, there are significant contrasts between this EELR and those of more extended radio galaxies.

First, the EELR of 9C J1503+4528 has a very sharp cut-off, rather than the more gradual decline in emission-line luminosity observed in the cases of other sources (e.g. Best et al. 2000; Inskip et al. 2002). Secondly, and of great significance, is the fact that the EELR is a factor of 10 larger than the size of the radio source. Unlike the case

² MP2DFUNFIT.PRO, part of Craig Markwardt's MPFIT non-linear least-squares curve-fitting package available via <http://astro.physics.wisc.edu/~craigm/idl/fitting.html>.

of the 3C and 6C sources, where the EELR is significantly affected by the passage of the radio source, here we have a case where the EELR could very well correspond to the limiting distance reached by the ionization front of an AGN, illuminating the region ahead of the radio source before it leaves the centre of the host galaxy.

This factor of ≈ 10 difference between radio source size and EELR size has a significant implication. In combination with the typical 0.1c source expansion speed, we find it highly plausible that the radio jets and AGN were triggered simultaneously, some $\sim 10^4$ yr prior to our observations.

6.3 Spatial distribution of star formation

The material lying at larger angles to the radio source axis (i.e. the regions probed by the perpendicular spectrum) is also of interest. The line emission along this axis is far less extensive (as expected for the regions lying on the edge of, or outside, the ionization cone). Its ionization state suggests a lower ionization parameter U (i.e. fewer and/or less energetic ionizing photons), and the weak line emission in these regions could be better explained as being due to ionization by young OB stars (H II regions). Our continuum modelling in the preceding section finds clear evidence for a YSP with an age of roughly 5 Myr, but does not allow us to pinpoint its location. However, the varying strength of the Balmer absorption features in the parallel and perpendicular spectra (Fig. 13) can give us some idea of how widely distributed this YSP might be.

This plot displays the $H\beta$ line profile for the parallel and perpendicular spectra, using increasing extraction widths of 1 to 5 arcsec, beyond which point increased aperture widths only lead to a deterioration of the S/N. In the perpendicular extracted spectra, Balmer absorption can be clearly observed, and becomes increasingly apparent as emission from regions at a greater distance from the AGN is included (i.e. for increased aperture widths). The line emission in the parallel spectra is roughly 50–100 per cent more luminous (see Appendix A), and the evidence for the same trend in the strength of the absorption component is less clear-cut. Whilst the line emission is strongest in the central regions of the host galaxy, the distribution of the line absorption must therefore be less well peaked, otherwise we would not observe a strengthening of the absorption component in wider apertures. This suggests that any YSP responsible for the Balmer absorption features is unlikely to lie predominantly within the central few pc. The pronounced appearance of the Balmer absorption in the perpendicular spectra suggests that the underlying YSP is not linked to the location of the growing radio source but rather, more widely distributed throughout the galaxy. Indeed, given the relative ages of YSP and radio source (5×10^6 and $\sim 10^4$ yr, respectively), a radio-source induced star formation scenario would seem unlikely.

It is interesting to contrast the derived properties of the YSP detected in the spectra of 9C J1503+4528 with those found for other sources. Recent work on the GPS radio source PKS1345+12 (Rodriguez et al., in preparation) finds evidence for a very similar YSP to that of 9C J1503+4528: an age of roughly 5 Myr (cf. $\sim 10^3$ yr for the radio source itself) and a mass of $\sim 10^6 M_\odot$. Imaging observations locate the YSPs in several super star clusters lying in the outer regions of the host galaxy, rather than in a nuclear starburst ring such as that observed for Cygnus A (e.g. Jackson, Tadhunter & Sparks 1998). A further point of interest is the finding that the hosts of high-luminosity AGN observed as part of the Sloan Digital Sky Survey (Kauffmann et al. 2003) typically have younger mean stellar ages than normal elliptical galaxies existing at the same

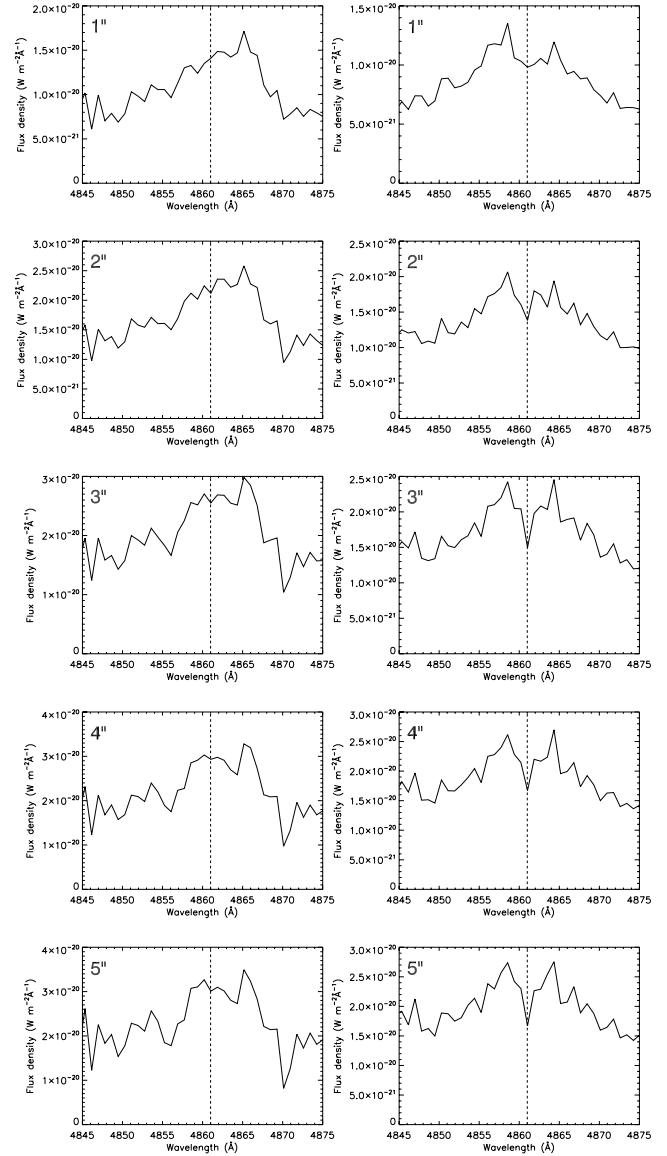


Figure 13. $H\beta$ line profiles for the parallel (left-hand panel) and perpendicular (right-hand panel) spectra, extracted using increasing aperture widths of (from top): 1, 2, 3, 4 and 5 arcsec. Balmer absorption is most clear in the perpendicular spectra, and generally becomes more pronounced with the inclusion of emission at greater distances from the AGN (i.e. the larger extraction apertures).

epoch. Additionally, the young stars are not preferentially located near the galaxy nuclei, but rather are spread out over a much wider scale. These properties are strikingly similar to those observed in 9C J1503+4528.

6.4 Gas kinematics

The observed changes in gas kinematics and dominant ionization mechanism in different regions of the EELR are not unexpected, given the expected location of radio source shocks, but do provide two key results: the clumpy ISM/IGM responsible for the EELRs observed around powerful radio sources does appear to be in place prior to the expansion of the radio source, and the kinematic properties of the extended regions are relatively undisturbed. Any

disturbance of the ISM/IGM due to an earlier interaction/merger has since settled down.

6.5 Spectroscopic evidence of a prior merger

If the triggering of radio source activity in this galaxy was due to a merger or interaction, the properties of any YSP formed as part of that process can shed light on the exact mechanisms and time-scales involved. Our continuum modelling of these spectra confirms the presence of a YSP, whose formation occurred several million years in advance of the onset of AGN activity. The mass of this YSP component is a low proportion of the total stellar mass of the host galaxy. Together with the quiescent kinematics in the outer regions of the EELR, and the morphological evidence described above, this suggests that *if* a merger/interaction was responsible for both the eventual triggering of the radio source activity and the formation of the YSP, it is most likely to have been fairly minor in nature.

7 CONCLUSIONS

9C J1503+4528 provides a detailed insight into many of the events associated with the triggering of a powerful radio source. Our observations have shown the following results.

(i) The relative sizes of EELR and radio source (the former being ≈ 10 times larger), together with typical compact radio source expansion velocities of $0.1c$, suggest that both the luminous quasar nucleus responsible for photoionizing the outer regions of the EELR, and the radio source itself, were triggered simultaneously.

(ii) The EELR ionization state is consistent with shocks being important on scales comparable to the radio source, with AGN photoionization within the ionization cones being the dominant mechanism further from the AGN, beyond the extent of the radio lobes.

(iii) The gas kinematics confirm that weak shocks could be important close to the radio source, and that the outer regions of the EELR are relatively undisturbed. The kinematics suggest that motion along the radio axis may be taking place, which may indicate the presence of outflows.

(iv) The blue continuum emission and Balmer absorption features suggest that recent star formation has occurred, but their spatial distribution suggests that it is not caused by the passage of the radio source.

(v) Modelling of the observed continuum emission gives a best-fitting age for the YSP of 5×10^6 yr, and a mass of roughly 0.06 per cent of the total stellar mass of the host galaxy.

(vi) Imaging observations of the host galaxy reveal a simple elliptical morphology; fitting of the K -band image suggests an effective radius of ~ 11.5 kpc.

(vii) A time-scale of events can be built up from these results: star formation is triggered by some event; in this case, any tidal interaction or merger must have been relatively minor. AGN/radio source activity follows the onset of star formation by several million years.

As a final consideration, we speculate on scenarios where the *lack* of evidence for a recent merger may be significant. As noted in Section 1, it is thought that the crucial ingredient for a radio-loud AGN is that the central black hole is rapidly spinning, a process which is most naturally explained by the coalescence of two smaller black holes (Wilson & Colbert 1995). Indeed, the cosmic evolution of AGN activity can be well described in terms of the expected halo merger rate (Mahmood, Devriendt & Silk 2005). However, the time-scales involved in the black hole coalescence after the halo

merger are extremely uncertain; while it is thought that supermassive black holes (SMBHs) can be brought within ~ 1 parsec of each other on a time-scale similar to the dynamical friction time-scale (e.g. Milosavljević & Merritt 2001), the mechanism for the final merger is unknown.

One possibility is that the formation of an accretion disc around one or both black holes provides the final necessary dynamical friction (e.g. Armitage & Natarajan 2005). If we assume that the central black hole of 9C J1503+4528 was spun up as part of the AGN triggering process, the apparent simultaneity of the triggering of both the optical AGN activity (requiring an accretion disc) and the radio jets (requiring the SMBH to be spinning rapidly) for 9C J1503+4528 would demand that the final SMBH coalescence must take place on a very short time-scale indeed. If this were not the case, and SMBH coalescence occurred more slowly, one would expect a period of radio-quiet AGN activity to precede the formation of the jets, which is ruled out in the case of this source by the relative sizes of the EELR and radio source.

Given the lack of evidence for a recent major merger, we therefore propose that the host galaxy of 9C J1503+4528 contained a rapidly spinning SMBH *prior* to the present period of AGN activity. In this case, the SMBH would have been spun up a significant time in the past, most likely during one or more mergers of smaller gas-rich discs which, in hierarchical clustering models, are the progenitors of ellipticals (e.g. Springel et al. 2005). This scenario has the advantage that radio-loud AGN may be ‘re-lit’ by relatively small gas injections from time to time, which provides a natural route for the necessary mechanical energy feedback in the intracluster medium (Best et al. 2005; Nipoti & Binney 2005). However, both 9C J1503+4528 and the well-studied cluster radio sources are observed at relatively recent cosmic time. The time-delay between the initial creation of a binary SMBH pair and the onset of jet activity may have been significantly shorter in the early Universe, where gas-rich mergers were significantly more prevalent. Such a scenario could be tested by investigating any time-lag between the evolution of the radio luminosity function at high redshift (which is still poorly constrained, e.g. Willott et al. 2001) and the halo merger rate as predicted by hierarchical clustering models.

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REFERENCES

- Armitage P. J., Natarajan P., 2005, *ApJ*, 634, 921
 Baldwin J. A., Phillips M. M., Terlevich R., 1981, *PASP*, 93, 5
 Best P. N., Röttgering H. J. A., Longair M. S., 2000, *MNRAS*, 311, 1
 Best P. N., Kauffmann G., Heckman T. M., Brinchmann J., Charlot S., Ivezić Ž., White S. D. M., 2005, *MNRAS*, 362, 25
 Binette L., Wilson A. S., Storchi-Bergmann T., 1996, *A&A*, 312, 365
 Blandford R. D., Znajek R. L., 1977, *MNRAS*, 179, 433
 Bolton R. C., Cotter G., Pearson T. J., Pooley G. G., Readhead A. C. S., Riley J. M., Waldrum E. M., 2003, *New Astron. Rev.*, 47, 367
 Bolton R. C. et al., 2004, *MNRAS*, 354, 485
 Bolton R. C., Chandler C. J., Cotter G., Pearson T. J., Pooley G. G., Readhead A. C. S., Riley J. M., Waldrum E. M., 2006, *MNRAS*, 367, 323
 Bremer M. N., Fabian A. C., Crawford C. S., 1997, *MNRAS*, 284, 213
 Brüggemann M., Kaiser C. R., 2002, *Nat*, 418, 301
 Bruzual G., Charlot S., 2003, *MNRAS*, 344, 1000
 Croton D. J. et al., 2006, *MNRAS*, 365, 11
 Dickson R., Tadhunter C., Shaw M., Clark N., Morganti R., 1995, *MNRAS*, 273, L29
 Dopita M. A., Sutherland R. S., 1996, *ApJS*, 102, 161
 Fabian A. C., Sanders J. S., Allen S. W., Crawford C. S., Iwasawa K., Johnstone R. M., Schmidt R. W., Taylor G. B., 2003, *MNRAS*, 344, L43
 Fanti C., Fanti R., Dallacasa D., Schilizzi R. T., Spencer R. E., Stanghellini C., 1995, *A&A*, 302, 317
 Groves B. A., Dopita M. A., Sutherland R. S., 2004a, *ApJS*, 153, 9
 Groves B. A., Dopita M. A., Sutherland R. S., 2004b, *ApJS*, 153, 75
 Heckman T. M., Smith E. P., Baum S. A., van Breugel W. J. M., Miley G. K., Illingworth G. D., Bothun G. D., Balick B., 1986, *ApJ*, 311, 526
 Howarth I. D., 1983, *MNRAS*, 203, 301
 Inskip K. J., Best P. N., Rawlings S., Longair M. S., Cotter G., Röttgering H. J. A., Eales S. A., 2002, *MNRAS*, 337, 1381
 Inskip K. J., Best P. N., Longair M. S., Röttgering H. J. A., 2005, *MNRAS*, 359, 1393
 Jackson N., Tadhunter C., Sparks W. B., 1998, *MNRAS*, 301, 131
 Johnston H. M., Hunstead R. W., Cotter G., Sadler E. M., 2005, *MNRAS*, 356, 515
 Kaiser C. R., Dennett-Thorpe J., Alexander P., 1997, *MNRAS*, 292, 723
 Kauffmann G. et al., 2003, *MNRAS*, 346, 1055
 Kewley L. J., Geller M. J., Jansen R. A., 2004, *AJ*, 127, 2002
 Mahmood A., Devriendt J. E. G., Silk J., 2005, *MNRAS*, 359, 1363
 Milosavljević M., Merritt D., 2001, *ApJ*, 563, 34
 Murgia M., 2003, *Publ. Astron. Soc. Aust.*, 20, 19
 Nipoti C., Binney J., 2005, *MNRAS*, 361, 428
 Oke J. B. et al., 1995, *PASP*, 107, 375
 Osterbrock D. E., 1989, *Astrophysics of Gaseous Nebulae and Active Galactic Nuclei*. University Science Books, Mill Valley, CA
 Owsianik I., Conway J. E., Polatidis A. G., 1998, *A&A*, 336, L37
 Polatidis A. G., Conway J. E., 2003, *PASA*, 20, 19
 Readhead A. C. S., Taylor G. B., Pearson T. J., Wilkinson P. N., 1996, *ApJ*, 460, 634
 Rola C. S., Terlevich E., Terlevich R. J., 1997, *MNRAS*, 289, 419
 Schlegel D. J., Finkbeiner D. P., Davis M., 1998, *ApJ*, 500, 525
 Sérsic J. L., 1968, *Atlas de Galaxias Australes*. Observatorio Astronomico, Córdoba
 Snellen I. A. G., Schilizzi R. T., Miley G. K., de Bruyn A. G., Bremer M. N., Röttgering H. J. A., 2000, *MNRAS*, 319, 445
 Springel V. et al., 2005, *Nat*, 435, 629
 Taylor G. B., Marr J. M., Pearson T. J., Readhead A. C. S., 2000, *ApJ*, 541, 112
 Venemans B. P. et al., 2005, *A&A*, 431, 793
 Waldrum E. M., Pooley G. G., Grainge K. J. B., Jones M. E., Saunders R. D. E., Scott P. F., Taylor A. C., 2003, *MNRAS*, 342, 915
 Willott C. J., Rawlings S., Blundell K. M., Lacy M., Eales S. A., 2001, *MNRAS*, 322, 536
 Wilson A. S., Colbert E. J. M., 1995, *ApJ*, 438, 62

APPENDIX A: EMISSION-LINE PROPERTIES OF 9C J1503+4528

Here, we present the results of our emission-line fitting for the parallel, perpendicular and combined two-dimensional frames. Fluxes relative to both $[\text{O II}] \ 3727 \text{ \AA}$ and $H\beta$ are tabulated in Table A1, as well as their equivalent widths. Although flux ratios quoted relative to $H\beta$ are more common, the presence of Balmer absorption lines reduces the usefulness of this ratio. The values in Table A1 have been corrected for galactic extinction, using values for the $H I$ column density of the Milky Way taken from the NASA Extragalactic Data base (NED), and the parametrized galactic extinction law of Howarth (1983).

Table A1. Spectroscopic properties of 9C1503+4528, determined from various spectra obtained perpendicular and parallel to the radio source. The spectra were extracted with various width apertures, so as to probe the properties of the central and exterior regions of the extended emission structures. The total integrated fluxes for the [O II] and H β emission lines are provided, given in units of $10^{-17} \text{ W m}^{-2}$. All other emission lines are quoted in terms of their flux ratios relative to both [O II] and H β within the extracted one-dimensional spectra. This has been done due to the fact that the levels of Balmer absorption vary with position throughout the galaxy, and therefore a comparison with the H β line flux, though popular in the literature, may not be the most appropriate option for this source. The errors on the H β and [O II] line fluxes are dominated by calibration errors, estimated to be $\lesssim 10$ per cent. Errors on the other lines and this calibration error are added in quadrature with the errors arising from photon statistics. All quoted values in this table are corrected for galactic extinction using the $E(B - V)$ for the Milky Way (Schlegel, Finkbeiner & Davis 1998) and the parametrized galactic extinction law of Howarth (1983).

Details of extracted spectrum		Combined, 2 arcsec	Parallel, 2 arcsec	Perpendicular, 2 arcsec	Parallel, 5 arcsec	Perpendicular, 5 arcsec
[O II] flux		191.5	227.5	158.6	248.0	180.2
H β flux		11.7	13.4	10.0	14.4	10.2
[O II] 3727	[O II] Flux ratio	100	100	100	100	100
	H β Flux ratio	1635 ± 177	1701 ± 191	1588 ± 195	1715 ± 231	1766 ± 290
	Equivalent width	182 Å	192 Å	170 Å	150 Å	123 Å
[Ne III] 3869	[O II] Flux ratio	8.5 ± 0.9	8.7 ± 1.0	8.1 ± 1.0	9.5 ± 1.1	11.1 ± 1.5
	H β Flux ratio	139 ± 16.2	148 ± 18.2	128 ± 17.9	162 ± 24.3	196 ± 36.7
	Equivalent width	17 Å	19 Å	15 Å	16 Å	15 Å
H ζ 3889	[O II] Flux ratio	0.7 ± 0.2	1.6 ± 0.3	1.1 ± 0.3	1.3 ± 0.5	< 0.3
	H β Flux ratio	12.1 ± 3.4	26.4 ± 5.8	16.7 ± 5.6	21.6 ± 7.9	< 5.2
	Equivalent width	1.3 Å	3.1 Å	1.9 Å	1.9 Å	< 0.2 Å
H ϵ + [Ne III] 3967	[O II] Flux ratio	0.6 ± 0.2	1.0 ± 0.2	< 0.3	0.6 ± 0.3	< 0.6
	H β Flux ratio	9.1 ± 3.1	16.4 ± 3.0	< 5.3	10.8 ± 4.8	< 10.3
	Equivalent width	1.0 Å	1.9 Å	< 0.3 Å	1.0 Å	< 0.4 Å
[S II] 4072	[O II] Flux ratio	2.7 ± 0.4	3.0 ± 0.4	2.7 ± 0.5	1.9 ± 0.4	2.0 ± 0.7
	H β Flux ratio	43.7 ± 6.5	50.0 ± 7.3	43.3 ± 9.1	31.7 ± 7.0	35.0 ± 13.5
	Equivalent width	4.1 Å	4.9 Å	3.9 Å	2.3 Å	2.1 Å
H δ 4102	[O II] Flux ratio	0.7 ± 0.3	0.9 ± 0.4	0.6 ± 0.4	0.8 ± 0.7	< 0.5
	H β Flux ratio	11.6 ± 4.8	14.4 ± 6.3	9.4 ± 6.0	13.6 ± 12.4	< 8.3
	Equivalent width	1.1 Å	1.5 Å	0.9 Å	1.1 Å	< 0.4 Å
H γ 4340	[O II] Flux ratio	3.3 ± 0.5	3.7 ± 0.5	2.8 ± 0.5	3.7 ± 0.6	1.5 ± 0.5
	H β Flux ratio	54.2 ± 7.7	63.0 ± 8.7	44.1 ± 8.5	62.2 ± 11.3	25.4 ± 10.1
	Equivalent width	5.7 Å	6.9 Å	4.3 Å	5.3 Å	1.7 Å
[O III] 4363	[O II] Flux ratio	1.2 ± 0.3	1.1 ± 0.3	0.8 ± 0.3	1.4 ± 0.4	0.7 ± 0.4
	H β Flux ratio	19.8 ± 5.1	19.1 ± 5.6	12.8 ± 4.8	23.2 ± 7.4	12.0 ± 6.7
	Equivalent width	2.1 Å	2.0 Å	1.2 Å	1.9 Å	0.8 Å
H β 4861	[O II] Flux ratio	6.1 ± 0.7	5.9 ± 0.7	6.3 ± 0.8	5.8 ± 0.8	5.7 ± 0.9
	H β Flux ratio	100	100	100	100	100
	Equivalent width	10 Å	10 Å	8.9 Å	7.7 Å	6.0 Å
[O III] 4959	[O II] Flux ratio	10.7 ± 1.1	9.9 ± 1.1	10.5 ± 1.2	13.5 ± 1.6	12.4 ± 1.6
	H β Flux ratio	175 ± 19.7	169 ± 20.2	167 ± 21.8	232 ± 33.5	219 ± 39.6
	Equivalent width	17 Å	17 Å	16 Å	18 Å	14 Å
[O III] 5007	[O II] Flux ratio	22.6 ± 2.4	23.7 ± 2.5	24.1 ± 2.7	30.5 ± 3.2	25.2 ± 3.3
	H β Flux ratio	370 ± 41.2	403 ± 46.3	383 ± 50.7	523 ± 72.9	446 ± 83.1
	Equivalent width	38 Å	43 Å	40 Å	41.8 Å	32 Å

APPENDIX B: BEST-FITTING YSP PARAMETERS IN THE PRESENCE OF A POWER LAW

As part of our modelling, we consider the effects of the inclusion of a power-law component on the best-fitting YSP parameters. (Table B1) The preferred age of the YSP remains unchanged at 5 Myr except for the strongest power-law contributions (>25 per cent). However, the reduced- χ^2 values indicate that models with a strong power-law contribution are not favoured.

Table B1. Changes to the best-fitting YSP model with the inclusion of a power-law component. Columns 1 and 2 give the percentage of the power-law component at a wavelength of ~ 3600 Å, and its spectral index, α (defined as $F_\lambda \propto \lambda^\alpha$). Column 3 gives the best-fitting old stellar population age in Gyr. The best-fitting YSP age (in Myr) and mass (as a percentage of the total stellar mass of the host galaxy) are given in columns 4 and 5 respectively. The reduced χ^2 for the best-fitting model is listed in column 6.

Power-law percentage at ~ 3600 Å (1)	α (2)	Old SP age (Gyr) (3)	YSP properties age (Myr) (4)	mass per cent (5)	Reduced χ^2 (6)
0.0	N/A	6.0	5	0.065	1.01943
1.0	-1	6.0	5	0.064	1.03379
1.0	-2	6.0	5	0.064	1.03269
1.0	-3	6.0	5	0.064	1.03236
2.5	-1	6.0	5	0.063	1.05603
2.5	-2	6.0	5	0.063	1.05338
2.5	-3	6.0	5	0.062	1.05281
5.0	-1	6.0	5	0.061	1.09503
5.0	-2	6.0	5	0.060	1.09005
5.0	-3	6.0	5	0.060	1.08985
10.0	-1	6.0	5	0.058	1.18089
10.0	-2	5.0	5	0.063	1.17021
10.0	-3	5.0	5	0.061	1.17095
15.0	-1	5.0	5	0.061	1.27639
15.0	-2	5.0	5	0.058	1.26151
15.0	-3	5.0	5	0.056	1.26901
25.0	-1	5.0	5	0.052	1.50695
25.0	-2	5.0	5	0.048	1.49007
25.0	-3	7.0	50	0.703	1.50599
50.0	-1	6.0	50	0.489	2.20266
50.0	-2	6.0	100	0.774	1.96529
50.0	-3	7.0	100	0.647	1.83815

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